

Measurement of collisions between rubidium atoms and optically dark rubidium ions in trapped mixtures

Seunghyun Lee,^{*} K. Ravi, and S. A. Rangwala[†]*Raman Research Institute, Sadashivanagar, Bangalore 560080, India*

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We measure the collision-rate coefficient between laser-cooled rubidium (Rb) atoms in a magneto-optical trap (MOT) and optically dark Rb^+ ions in an overlapping Paul trap. In such a mixture, the ions are created from the MOT atoms and allowed to accumulate in the ion trap, which results in a significant reduction in the number of steady-state MOT atoms. A theoretical rate-equation model is developed to describe the evolution of the MOT atom number due to ionization and ion-atom collisions, and we derive an expression for the ion-atom collision-rate coefficient. The loss of MOT atoms is studied systematically by sequentially switching on the various mechanisms in the experiment. Combining the measurements with the model allows the direct determination of the ion-atom collision-rate coefficient. Finally, the scope of the experimental technique developed here is discussed.

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I. INTRODUCTION

Trapped ion and atom mixtures enable us to investigate exciting physics ranging from a single trapped ion inside a Bose-Einstein condensate (BEC) [1,2] and atoms in a magneto-optical trap (MOT) superimposed with laser-cooled ions [3–5] to mixtures where the ions are only cooled in collision with the atoms [6–8]. Elastic and charge-exchange processes in these mixed, trapped system have been demonstrated and measured [3,8,9]. Even before the existence of such mixed traps, a variety of rate coefficients and cross sections for several ion-atom processes have been measured [10]. The most prevalent method for the determination of the ion-neutral collision rate using trapped ions measures the loss of ions from the trap [10,11]. For optically dark ions, this loss is measured by sweeping the ion resonance frequency or by extracting the trapped ions to charge particle detectors [10]. However, in mixed ion-atom trapping experiments in which the ions are optically dark, the cold trapped atoms interacting with ions can be directly measured *in situ* by atomic fluorescence without affecting the trapped ions. In this article, we develop a method to measure ion-atom collision-rate coefficients in such a hybrid system using atom fluorescence as a probe.

In the experiment reported here, we trap the optically dark rubidium ($^{85}\text{Rb}^+$) ions derived from the laser-cooled ^{85}Rb atoms in a MOT. The ions are created and accumulated at the center of the ion trap, which overlaps with the MOT. Such a system evolves into a mixture with constant numbers of trapped atoms and ions in the steady state. The ionization of the MOT atoms and the subsequent trapped-ion interaction with the atoms results in a depletion of the atoms from the MOT. The consequent drop in MOT fluorescence is utilized to develop the framework for the measurement of the collision-rate coefficient for ion-atom interactions, as described below.

In what follows, we construct the method for the determination of the collision-rate coefficient. We first present a brief description of the experimental setup, followed by a simple,

analytical rate-equation model, where we systematically consider loading and loss processes of the MOT atoms. Here we model in a calibrated manner the change in the trapped-atom number when each ion-related process is turned on. This culminates with the simultaneous trapping of ions and atoms in steady state. The experimental data are then presented in a sequence consistent with the processes modeled. Combining the measurements with the derived analytical expressions allows the determination of the rate of atom loss due to ion-atom interactions, and hence the collision-rate coefficient for ions and atoms. We end with a discussion of some features of such measurements.

II. EXPERIMENTAL SETUP

The experimental arrangement (Fig. 1), techniques, and characterization have been described in detail in previous work [12]. Here we include a brief description for making the experimental discussion comprehensible. Rb-atom vapor is created from a dispenser and ^{85}Rb atoms are laser cooled and trapped in a six-beam MOT of rms size 0.35 mm and temperature $166 (\pm 28) \mu\text{K}$. The cooling laser is red detuned by $\delta_c = 2\pi \times 10$ MHz with respect to the $5S_{1/2}(F=3) \leftrightarrow 5P_{3/2}(F'=4)$ transition, and the laser power is 3.5 mW per beam. A magnetic gradient field of 12 Gauss/cm is utilized for the MOT. The intersection point of the MOT laser-intensity maxima and magnetic-field zero coincides with the center of a linear Paul trap, built around the capture volume of the MOT. This ensures that the trapped ions and laser-cooled atoms are optimally overlapped [6,12]. Ions are created from cold atoms using a blue-light source (BLS) which along with the cooling light results in two-photon ionization of the Rb MOT atoms. The ion trap is made up of four parallel rods in the quadrupole configuration, where a time-varying radio-frequency (rf) voltage is applied to one diagonal set, and an opposite phase rf voltage is applied to the complementary diagonal set. This generates a two-dimensional quadrupole field for the radial trapping of ions. Two end-cap ring electrodes at a positive dc voltage confine the positive ions along the trap axis. Atoms are detected by observing the fluorescence from the MOT on a photomultiplier tube (PMT) and imaged

^{*}lee@rri.res.in[†]sarangwala@rri.res.in

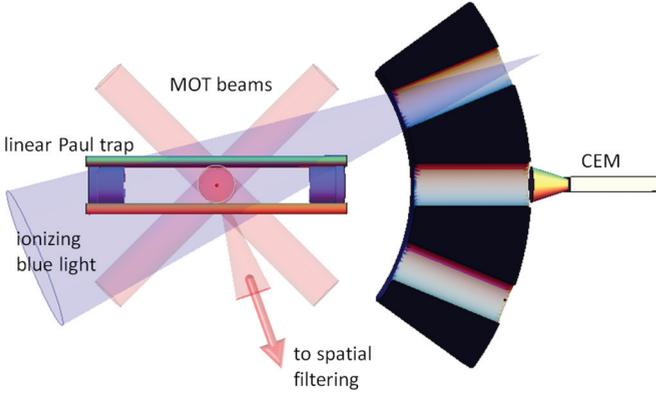


FIG. 1. (Color online) Schematic of experimental setup. The MOT and the linear Paul trap are overlapped as shown, with the cooling lasers illustrated in red, and the MOT located at the intersection point represented by the red sphere. The ionizing blue light is illustrated, without its source. The channel electron multiplier (CEM) is used to detect the trapped ions by appropriately switching the voltage on the hollow end-cap electrode closer to the CEM. The direction for the MOT fluorescence detection is indicated, without illustrating the spatial-filtering arrangement.

using two CCD cameras. The MOT fluorescence is measured by the PMT, where a spatial-filtering arrangement is used to reduce background-light signal. The ions are optically dark and are therefore detected either indirectly by a change in atomic fluorescence of the MOT or destructively by extraction onto a channel electron multiplier (CEM) and measuring the ion counts, either by pulse counting or by measurement of the proportional extracted ion signal. The ion extraction to the CEM is affected by switching the ion trap end-cap electrode from +80 V to -5 V while the rf field is still on.

III. MODEL FOR LOSS OF TRAPPED ATOMS

A. MOT rate equation

We now present the rate-equation framework which describes the evolution of the ion atom system as we add the various ingredients of the above experiment one at a time. The total number of atoms in the MOT (N_{MOT}) at any given time is $N_{\text{MOT}} = N_g + N_e$, where N_g and N_e are the atom numbers in the ground and excited states, respectively. The time dependence of N_{MOT} can be written as

$$\frac{dN_{\text{MOT}}}{dt} = L - \gamma_{ml}N_{\text{MOT}}, \quad (1)$$

with L being the loading rate of the atoms from the background vapor into the MOT, and γ_{ml} being the loss rate of the MOT atoms, which is primarily collisional [13]. The time-dependent solution of Eq. (1) with initial condition $N_{\text{MOT}}(0) = 0$ is

$$N_{\text{MOT}}(t) = N_0(1 - e^{-\gamma_{ml}t}), \quad N_0 = \frac{L}{\gamma_{ml}}, \quad (2)$$

where N_0 is the steady-state atom number in the MOT in the limit of large time.

B. Fraction of excited atoms

The fraction of atoms in the excited state f_e is determined from the photon scattering rate from an atom and can be written as [14]

$$f_e = \frac{N_e}{N_g + N_e} = \frac{(\Omega/\Gamma)^2}{1 + 4(\delta_c/\Gamma)^2 + 2(\Omega/\Gamma)^2}, \quad (3)$$

where $\Omega = 2\pi \times 16.7$ MHz is the Rabi frequency and $\Gamma = 2\pi \times 6$ MHz is the decay rate of the excited atomic state.

C. Photoionization of atoms

The ionization of atoms in our experiment is accomplished by excitation of the atoms in the $5p_{3/2}$ state to the ionization continuum using the BLS. The central wavelength the BLS is at $\lambda_{pi} = 456$ nm, which is sufficient to ionize the excited Rb atoms [15]. The intensity of the BLS light incident on the MOT is denoted by I_{pi} . For photoionization energy $E_{pi} = hc/\lambda_{pi}$, the ionization cross section is denoted by σ_{pi} . The photon energy which is in excess of the ionization energy is carried away by the ejected electron, with the resulting ion suffering the recoil.

The loading rate of the MOT remains unchanged when the BLS is switched on for photoionization, as the population of the background atoms remains unchanged with operation of the BLS. The maximum intensity of the BLS used is 1.33 mW/cm². The rate equation for the MOT, when the BLS is ON, is modified to

$$\frac{dN_{\text{MOT}}}{dt} = L - \gamma_{ml}N_{\text{MOT}} - \gamma_{pi}N_{\text{MOT}}, \quad (4)$$

where a new loss rate γ_{pi} adds to the losses inherent to the operation of the isolated MOT to give a total loss rate of $\gamma_t = \gamma_{ml} + \gamma_{pi}$ [16,17]. Here we assume that since the ion trap is switched off during photoionization, the ions and electrons created from the MOT atoms immediately leave the system and play no further role in determining the number of trapped atoms. The ionization rate of the excited Rb atoms by radiation from the BLS can be written as [16,17]

$$\gamma_{pi} = \left(\frac{\sigma_{pi} f_e}{E_{pi}} \right) I_{pi} = \zeta I_{pi}, \quad (5)$$

where $\zeta = (\frac{\sigma_{pi} \lambda_{pi} f_e}{hc})$. The corresponding time dependence of N_{MOT} with initial condition $N_{\text{MOT}}(0) = 0$ and the inclusion of the ionizing process can then be written as

$$N_{\text{MOT}}(t) = \frac{L}{\gamma_t}(1 - e^{-\gamma_t t}). \quad (6)$$

D. Ion-atom interaction loss

Once the ion trap is switched ON, then a fraction of the created ions load into the ion trap which overlaps the MOT. Since the ion and atom traps are optimally overlapped, the ions created from the MOT have negligible velocity upon creation. These ions gain in kinetic energy as they evolve in the field of the ion trap. As the number of ions loaded in the ion trap grows, the spatial distribution of the trapped ions increases outward from the trap centers. Since the MOT volume is much smaller than the volume of the trapped-ion spatial distribution,

then long before the total number of trapped ions equilibrates, the number of ions at any instant overlapped with the MOT stabilizes. Therefore the collision rate between the trapped ions and the MOT atoms also stabilizes rapidly, allowing the definition of a time-independent binary ion-atom interaction rate, γ_{ia} for the constituents of the ion-atom mixture. Since atoms are weakly trapped with respect to ions, they can gain sufficient energy in collision with trapped ions to exit the MOT. This is a new loss channel for the atoms trapped in a MOT, which can be written as

$$\frac{dN_{\text{MOT}}}{dt} = L - (\gamma_{ml} + \gamma_{pi} + \gamma_{ia})N_{\text{MOT}}. \quad (7)$$

The resulting time evolution for N_{MOT} with initial condition $N_{\text{MOT}}(0) = 0$ therefore becomes

$$N_{\text{MOT}}(t) = \frac{L}{\gamma_{\text{tot}}}(1 - e^{-\gamma_{\text{tot}}t}), \quad (8)$$

where $\gamma_{\text{tot}} = \gamma_t + \gamma_{ia}$ is the total loss rate of MOT atoms when the ion and atom traps are operated simultaneously as described in the experimental section above. We note that, while the ion density overlapping with the MOT volume stabilizes rapidly, the number and the velocity distribution of the trapped ions evolve continuously until the ion trap is filled to its limit. Thus the rate of ion loss from the ion trap and the evolution of N_{MOT} are linked.

E. Number of trapped ions

The ion trap is loaded to the minimum of the secular trap potential, by photoionization of the overlapping MOT atoms. The number of trapped ions as a function of time, for a given photoionizing intensity I_{pi} , can be described as

$$N_I(t) = N_{\text{MOT}}\zeta \frac{I_{pi}}{\gamma}(1 - e^{-\gamma t}), \quad (9)$$

where γ is the ion-trap loss rate. Since ions have much higher velocities than the MOT capture velocity for the atoms, ion-atom collisions will lead to loss of atoms from the MOT. This allows the number of MOT atoms lost due to trapped ions, $N_{\text{MOT}}^{\text{loss}}$, to be written as

$$N_{\text{MOT}}^{\text{loss}}(t) = N'(1 - e^{-\gamma_{ia}t}), \quad (10)$$

where N' is the saturation number of atom loss due to ion-atom interaction, and γ_{ia} is the loss rate of the MOT atoms due to ion-atom interactions. The above form of the atom loss will be validated later [Eqs. (21) and (22) and Fig. 2], when the experiment and results are presented. The atom-loss rate from the MOT, due to ion-atom interactions is proportional to the number of trapped ions [$N_{\text{MOT}}^{\text{loss}}(t) \propto N_I(t)$], and so from Eqs. (9) and (10) we conclude that $\gamma \equiv \gamma_{ia}$. Therefore, number of trapped ions as a function of t is expressed as

$$N_I(t) = N_{\text{MOT}}\zeta \frac{I_{pi}}{\gamma_{ia}}(1 - e^{-\gamma_{ia}t}). \quad (11)$$

The number of trapped ions also depends on I_{pi} but cannot increase indefinitely with I_{pi} due to finite trap depth and volume. Therefore an intensity-loss coefficient of trapped ions must be introduced. In the same spirit as the above discussion,

to construct a general form of N_I , we write

$$\frac{dN_I(t, I_{pi})}{dI_{pi}} = N_{\text{MOT}}\zeta \frac{1}{\gamma_{ia}}(1 - e^{-\gamma_{ia}t}) - \kappa N_I(t, I_{pi}), \quad (12)$$

where κ is the intensity-loss coefficient and has units of inverse intensity. The solution of the above equation is then

$$N_I(t, I_{pi}) = N_{\text{MOT}}\zeta \frac{1}{\gamma_{ia}\kappa}(1 - e^{-\gamma_{ia}t})(1 - e^{-\kappa I_{pi}}). \quad (13)$$

The equation can be fully converted into a time-dependent function of trapped ions when $I_{pi} \rightarrow \infty$ and also fully converted into an I_{pi} -dependent function of trapped ions when $t \rightarrow \infty$. The number of ions trapped as a function of I_{pi} for $t \rightarrow \infty$ is written as

$$N_I(I_{pi}) = N_I^0(1 - e^{-\kappa I_{pi}}). \quad (14)$$

The maximum trapped ion number, N_I^0 when $t \rightarrow \infty$ and $I_{pi} \rightarrow \infty$ is expressed as

$$N_I^0 = N_{\text{MOT}}\zeta \frac{1}{\gamma_{ia}\kappa}. \quad (15)$$

Since Eqs. (14) and (15) apply to the ion system in the limit $t \rightarrow \infty$, it is therefore reasonable to expect an average motional energy $\langle E_I \rangle$ for each ion.

F. Determination of ion-atom collision rate

The ions trapped as described above are assumed to have a speed distribution $f(v)$, consistent with $\langle E_I \rangle$. The laser-cooled atoms on the other hand are expected to exhibit a temperature $T_A \approx 100 \mu\text{K}$. In this situation $\langle E_I \rangle \gg k_B T_A$ and therefore all of the velocity of the ion-atom collision can be assumed to be possessed by the ion, in the laboratory frame of reference (LFor). Similarly, as is shown later, $N_{\text{MOT}} \gg N_I^0$ because the ion density ρ_I is far less than the atom density ρ_A in such mixtures despite the ion trap volume $V_{\text{IT}} \gg V_{\text{MOT}}$, which is the MOT volume in our experiment.

The total ion-atom collision cross section σ_{tot} is energy ($E \propto v^2$) dependent and is the sum of the elastic, σ_{el} , and resonant charge exchange, σ_{cx} . $\sigma_{\text{el}} \propto 1/E^{1/3}$ over all the energies and $\sigma_{\text{cx}} \propto 1/E^{1/2}$ in the low-energy regime and $\sigma_{\text{cx}} \propto (a \ln E - b)^2$ at high collision energies (a and b are constants) [18]. Given $f(v)$ for the ions, the determination of σ_{tot} rests on assumptions made for $f(v)$. In this case the experimentally meaningful quantity to consider is the ion-atom rate coefficient

$$k_{ia} = \int_v \sigma_{\text{tot}} v f(v) dv \equiv \langle \sigma_{\text{tot}} v \rangle, \quad (16)$$

which represents the velocity-averaged product of $\sigma_{\text{tot}} v$. A single ion with velocity v_0 and a corresponding ion-atom cross section σ_{tot}^0 will collide with MOT atoms of density ρ_A at a rate given by

$$z_0 = \sigma_{\text{tot}}^0 v_0 \rho_A. \quad (17)$$

For the $N_{I;M}$ ions that overlap the MOT at any given time, with a speed distribution $f(v)$, the total ion-atom collision rate then becomes

$$z N_{I;M} = \langle \sigma_{\text{tot}} v \rangle \rho_A N_{I;M} \equiv k_{ia} \rho_A N_{I;M}. \quad (18)$$

For a given value of I_{pi} , the number of ions in the ion-trap volume V_{IT} , overlapping the MOT volume V_{MOT} , is given by

$$N_{I;M} = N_I^0 (1 - e^{-\kappa I_{pi}}) \frac{V_{MOT}}{V_{IT}}. \quad (19)$$

Since ion velocities are large, and the MOT capture velocity for the atoms much smaller, ion-atom collisions cause the MOT atoms to eject, allowing us to equate the total ion-atom collision rate in Eq. (18) to $\gamma_{ia} N_{MOT}$, the atom-loss rate due to ion-atom collisions. After substituting Eq. (19) into Eq. (18), the resulting expression for γ_{ia} then becomes

$$\gamma_{ia} = \frac{N_I^0 k_{ia}}{V_{IT}} (1 - e^{-\kappa I_{pi}}), \quad (20)$$

which defines the relation between the experimentally measurable γ_{ia} and the rate coefficient for ion-atom collisions. In the experimental results that follow, we demonstrate that the above rate-equation formulation describes adequately the dependence of the various loss rates with I_{pi} and the BLS intensity.

IV. RESULTS

A. Experimental sequence

We now describe an experiment that allows us to validate the above rate-equation analysis for ion-atom interaction. The basic time sequence instrumental for results below is shown in Fig. 2. Here a MOT containing $\approx 1.80 (\pm 0.06) \times 10^6$ atoms is loaded to saturation in 40 s (T_1 to T_2), and the change in the number of atoms in the MOT is measured under two conditions from T_2 to T_3 at different intensities of the ionizing light (BLS). When the BLS ionizes the cooled ^{85}Rb atoms from the MOT and these ions and electrons leave the MOT volume immediately because there is no confinement for them. This results in a depletion of the steady-state atom number in the MOT and therefore allows the determination of the ionization rate for the excited atoms. Using Eq. (2), data from T_1 to T_2 are fit to determine γ_{ml} . The solution for Eq. (4) with appropriate initial conditions, $N(t=0) = N_0$ and $N(t \rightarrow \infty) = N_0 - N_t$, is

$$N_{MOT}(t) = N_0 - N_t (1 - e^{-\gamma_t t}), \quad (21)$$

where N_t is the atom loss due to γ_t during photoionization. The quantity γ_t is determined by fitting data from T_2 to T_3 in Fig. 2(a) with Eq. (21). In Fig. 2(b) we switch ON both the ionizing light (BLS) and the ion trap simultaneously. In this case the ionization process is unaffected, but some fraction of the ions created reside within the ion trap, which is well overlapped with the atom trap. Here we observe that the number of trapped atoms reduces significantly over the ionization loss. This loss is attributed to ion-atom interactions. The solution for Eq. (7) is obtained in a similar way as for the photoionization case and is

$$N_{MOT} = N_0 - N_{tot} (1 - e^{-\gamma_{tot} t}), \quad (22)$$

where N_{tot} is the atom loss due to γ_{tot} when both atoms and ions are simultaneously trapped.

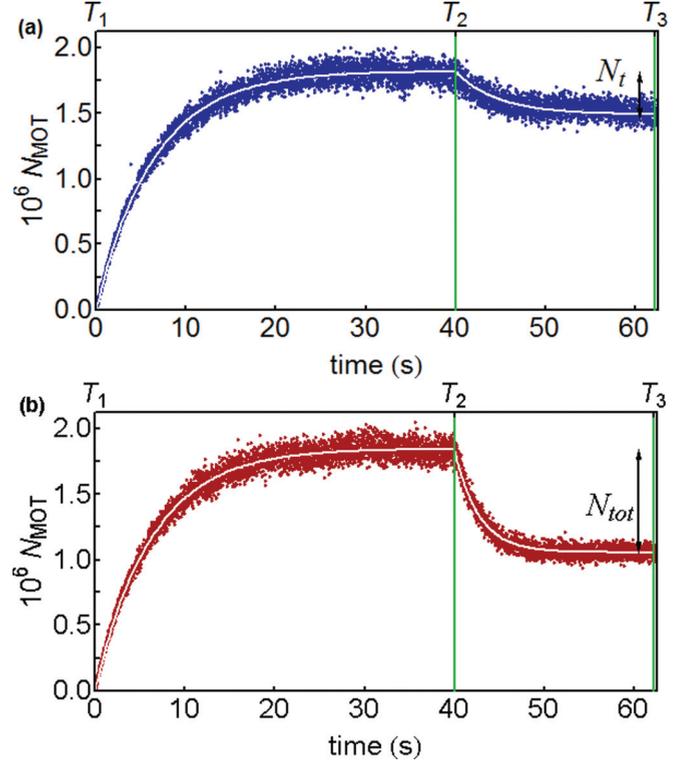


FIG. 2. (Color online) Temporal sequence for experiment. In both panels, first the MOT is loaded for 40 s (from T_1 and T_2) to saturation, where it contains $\approx 1.8 \times 10^6$ atoms. Beyond T_2 , the photoionization light is switched on. Panels (a) and (b) illustrate the evolution of the trapped atoms when the ion trap is OFF and with ion trap ON, respectively. The number of trapped atoms reduces in both cases. However, when the ions and the atoms are simultaneously trapped, the atom loss is much more significant. The plots in (a) and (b) correspond to BLS intensity $I_{pi} = 1.33 \text{ mW/cm}^2$. The various loss rates discussed in the text are determined from such data as a function of photoionizing intensity. N_t and N_{tot} are the atom losses due to γ_t and γ_{tot} . The lines show the fits to the data according to Eqs. (21) and (22), respectively.

B. Ionization without trapping

When the BLS is switched ON at T_2 , and the ion trap is not operational, and the loss rate of atoms from the MOT increases due to ionization of the atoms, which is shown in Fig. 3. The measured loss-rate coefficient $\gamma_t = \gamma_{ml} + \gamma_{pi}$ can be written using Eq. (5) as

$$\gamma_t = \gamma_{ml} + \zeta I_{pi} \quad (23)$$

From Fig. 3 it is clear that γ_t has a linear dependence on the intensity of the BLS, I_{pi} . This is because, although the ionization process is a two-photon process, the operation of the MOT ensures that a constant fraction of atoms are present in the excited state, making the ionization an effective single-photon process. The slope of γ_t , ζ is determined to be $0.00225 (\pm 0.00022) \text{ m}^2/\text{J}$ from the data. The value of γ_t for $I_{pi} = 0$ is $\gamma_{ml} = 0.144 (\pm 0.011) \text{ s}^{-1}$.

C. Ionization and ion trapping

For the case when ions are held in the ion trap after T_2 in Fig. 2(b), the loss-rate coefficient γ_{tot} varies with a

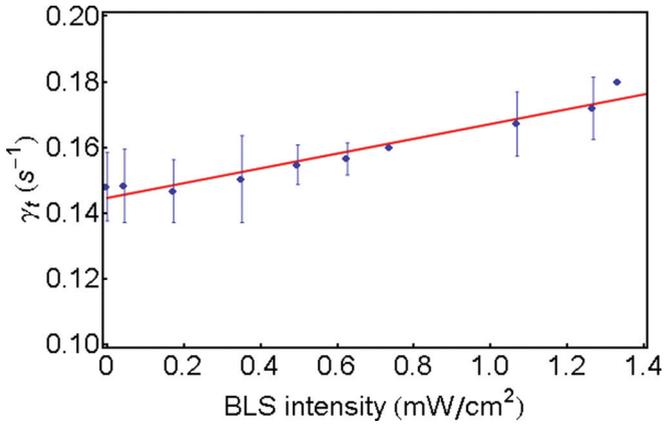


FIG. 3. (Color online) Loss rate of atoms from MOT as a function of photoionization light intensity, I_{pi} , without the ion trap in Fig. 2(a). As I_{pi} increases, the loss of atoms from the MOT increases as does the rate of loss γ_{pi} . Since this atom loss is happening on top of the normal loss from the MOT, the experimentally measured quantity here is γ_t . As can be seen from the text, a linear relation is expected between γ_t and I_{pi} , which is measured with a slope of $0.00225 (\pm 0.00022) \text{ m}^2/\text{J}$ and the γ_t -axis intercept of $\approx 0.144 \text{ s}^{-1}$ which is the MOT loss rate without the ionizing light (BLS). Error on the slope from the fit is calculated from confidence-level analysis.

nonlinear behavior as shown in Fig. 4. The change in γ_{tot} is fit by a combination of the linear expression in Eq. (23) and the exponential form of Eq. (20). To study the loss rate due to ion-atom interaction, γ_{ia} is plotted separately, as $\gamma_{ia} = \gamma_{tot} - \gamma_t$, in Fig. 5. The exponential nature of γ_{ia} as a function of I_{pi} is evident, thus validating the rate equation for the ion-atom interaction. From the fit to the data, Eq. (20) becomes $\gamma_{ia} = 0.218(1 - \exp[-0.126I_{pi}])$. Identifying $\kappa = 0.126 (\pm 0.031)$ in inverse intensity units, i.e., m^2/W and

$$\frac{N_I^0 k_{ia}}{V_{IT}} = 0.218 (\pm 0.029), \quad (24)$$

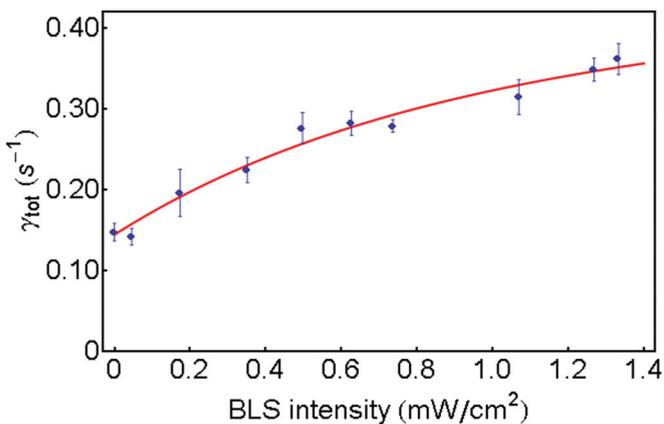


FIG. 4. (Color online) Loss rate of atoms from MOT as a function of ionization light intensity I_{pi} , with the ion trap in Fig. 2(b). As I_{pi} increases, the loss of atoms occurs on top of the normal loss from the MOT and the ionization loss, and so the experimentally measured quantity here is γ_{tot} . The additional loss of atoms over that in Fig. 3 is attributed to the ion-atom interactions for which the loss rate is γ_{ia} . γ_{tot} has exponential and linear components in I_{pi} as is observed.

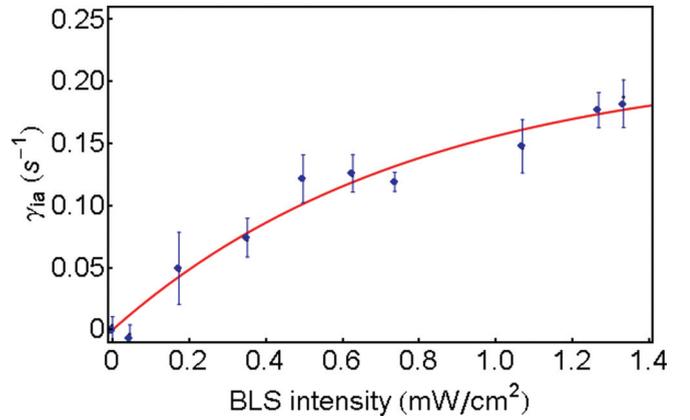


FIG. 5. (Color online) Subtraction of data in Fig. 3 from those in Fig. 4 is plotted to isolate the ion-atom interaction term. This loss rate of the MOT atoms due to ion-atom interaction, γ_{ia} , is plotted against BLS intensity, I_{pi} . γ_{ia} is fit with a single-exponential function of the form of Eq. (20) and the coefficients of the fit are utilized to characterize the interaction term, as discussed in the text.

with units s^{-1} , we can constrain the product of the trapped ion density and the ion-atom collision-rate coefficient.

D. Direct ion detection

Alongside the measurements on the MOT discussed above, the trapped ions though optically dark in the present experiment, can be measured using a CEM as described in the experimental arrangement. At the end of each ion-atom experiment, the trapped ions are extracted into the CEM by switching the end-cap voltage appropriately. However, since this experiment fills the ion trap to its capacity, severe pile-up results due to overlapping arrival times of the ions onto the detector. For a CEM cone voltage of -2100 V , the extracted ions from the trap as a function of I_{pi} is illustrated in Fig. 6. A single-exponential dependence of the ion numbers detected with BLS intensity can be expected and a dependence of that nature is seen in the data. However, there is detector

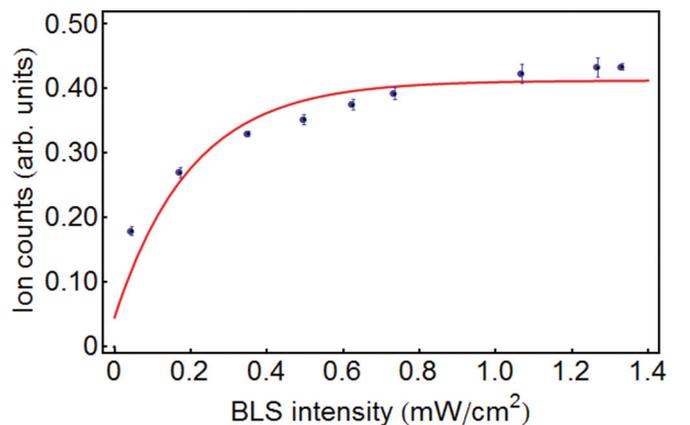


FIG. 6. (Color online) Change in detected ion counts vs BLS intensity. Since the number of trapped ions is very large, severe pileup of the ion signal results and the ion signal is seen to saturate faster than measured from the atom signal. The fit is generated using Eq. (14).

saturation to contend with and the present measurement of ion numbers is limited by the detection. The direct measurement of the CEM signal is therefore presented for the sake of experimental completeness rather than with the motive of utilizing it to determine the ion-atom interaction rates. For large ion numbers this is likely to remain a significant problem with present-day ion detectors. Alternative schemes for ion detections by electrode pickup, Faraday cup detection, etc. [19], may be useful for this problem and should be explored.

E. Collision-rate-coefficient determination

We now demonstrate the use of Eq. (24) to arrive at the ion-atom collision rate in the present experiment. The values of all the quantities used in the determination of the collision rate, with the standard deviation error in the accompanying parenthesis are provided in the text below. The product of the ion density and k_{ia} is determined from the fit to Eq. (20) as $0.218 (\pm 0.029) \text{ s}^{-1}$, so only the determination of the number of trapped ions and the volume of trapping for the ions is required for computation of the collision-rate coefficient. Keeping the compromised CEM data in mind, $\kappa = 0.126 (\pm 0.031) \text{ m}^2/\text{W}$ is determined from the experimental data for atom loss in Fig. 5, rather than from the data in Fig. 6. Using Eq. (15), the maximum number of ions which can be accumulated in the ion trap volume is calculated to be $N_I^0 = 1.48 (\pm 0.43) \times 10^5$.

The volume within which the ions are trapped, V_{IT} , is determined for the present trap parameters by a combination of trap loss and trap secular frequency measurements. A Monte Carlo analysis of the time for the loss of ions from the ion trap [6] gives the secular trap depth for trapped ions to be $\approx 0.8 \text{ eV}$. The secular frequencies for the ion trap are $\omega_x = \omega_y = 2\pi \times 135 \text{ kHz}$, and $\omega_z = 2\pi \times 27 \text{ kHz}$, which allow us to calculate the trap extent in each direction from the relation $m_I \omega^2 r^2 / 2 = 0.8 \text{ eV}$, where $r \in x_0, y_0, z_0$ is the extremal displacement in each direction and m_I is the mass of the ion. The trapping volume defined by these dimensions for a single ion then is $V_{IT} = 8.33 (\pm 0.83) \times 10^{-8} \text{ m}^3$.

The collision rate is determined by substituting the mean values of the quantities above in Eq. (24), which gives

$$k_{ia} = (1.23 \pm 0.42) \times 10^{-13} \text{ m}^3/\text{s}. \quad (25)$$

The standard deviation error is $\sigma_{k_{ia}} \equiv (\sum_i \sigma_i^2)^{1/2} = 0.42 \times 10^{-13} \text{ m}^3/\text{s}$, where the σ_i 's are the individual errors for each contributing parameter, shown in Table I. This demonstrates the technique's ability to arrive at the rate coefficients for ion-atom processes, even when the ions are not directly detected. The above rate coefficient is determined when the fraction of

atoms in the excited state for $f_e \approx 0.28$, from Eq. (3). The measured rate coefficient incorporates the elastic and resonant charge exchange collision from the ground and excited states of the atoms.

F. Theoretical estimate of collision-rate coefficient

To estimate the collision-rate coefficient in Eq. (16) we adopt the analytical form of total cross section as a function of collision velocity from Refs. [18,20] by using the ground-state C_4 value for Rb atoms, $5.26 \times 10^{-39} \text{ Cm}^2/\text{V}$ [14]. The speed distribution of trapped ions is determined using Monte Carlo analysis and molecular dynamics simulations [6]. Molecular dynamics simulation shows that the ion trap can stably trap ions with a Maxwell Boltzmann (MB) distribution of temperature $\approx 1000 \text{ K}$. In a Monte Carlo analysis, a noninteracting distribution of ions are evolved in the absence of cold atoms in the ion trap potential to estimate the maximum trappable secular energy, which is found to be $\approx 0.8 \text{ eV}$, corresponding to speed of 1360 m/s . This value matches the tail of the MB distribution of ions in the trap. Since $V_{IT} \gg V_{MOT}$ and the two traps are well centered in the overlap region, the micromotion velocity is small. Therefore, to estimate the speed distribution of the ions, only considering the secular velocity is a good approximation. By performing the integration in Eq. (16) with these quantities, we calculate the rate coefficient as $9.4 \times 10^{-14} \text{ m}^3/\text{s}$, when all the Rb atoms are in the ground state. Because atoms in the MOT are constantly pumped to the excited state by the cooling laser, ions invariably collide with a fraction of the MOT atoms, which are in the excited-state atoms. In the experiment, we have 28% of the atom population in the excited state ($f_e = 0.28$). The scalar polarizability for a Rb atom in the excited $5P_{3/2}$ state is $14.15 \times 10^{-39} \text{ Cm}^2/\text{V}$ [14], which is larger than that in the ground state. Incorporating this fraction of atoms in the excited state into the calculation, we compute the rate coefficient as $k_{ia} = 1.12 \times 10^{-13} \text{ m}^3/\text{s}$, which agrees well with the experimental value.

V. DISCUSSION

The above technique allows the detection of collisional processes between trapped ions and atoms. In the present case, this is successful even for optically dark ions. At its core it relies on the measured changes in the fluorescence of the MOT atoms in the presence of the ions. The present work relies on two important and reasonable premises, which need to be emphasized.

TABLE I. Key values and errors for the quantities used in the determination of k_{ia} .

Parameter	Value (\pm error)	Units	Method
ζ	$0.00225 (\pm 0.00022)$	m^2/J	Fit from experimental data in Fig. 3
N_{MOT}	$1.80 (\pm 0.06) \times 10^6$	Number	Fluorescence measurement
V_{IT}	$8.33 (\pm 0.83) \times 10^{-8}$	m^3	Derived from secular frequencies and trap depth
N_I^0	$1.48 (\pm 0.43) \times 10^5$	Number	See Eq. (15)
$\gamma_{ia}(I_{pi} \rightarrow \infty)$	$0.218 (\pm 0.029)$	s^{-1}	Fit from experimental data in Fig. 5
κ	$0.126 (\pm 0.031)$	m^2/W	Fit from experimental data in Fig. 5
k_{ia}	$1.23 (\pm 0.42) \times 10^{-13}$	m^3/s	Calculated from various quantities given above

The first is that ion-atom collisions are two-body processes. This is reasonable because of several conditions that exist in the experiment. The atoms in the MOT are noninteracting to a good approximation. Because the ions are hot, they have large velocities and in this regime the binary ion-atom cross section is small. This is enough to make the presence of another ion in the vicinity of the collision, such that it affects the details of the ion-atom collision, highly improbable. It should be noted, however, that the experiment and the model are in good agreement, and since the model only incorporates two-particle ion-atom collisions, it is fair to argue that any more complex process is absorbed within the errors quoted in the present result.

The second premise is that ion-atom collisions knock out MOT atoms. In the experiment, the only cooling mechanism for the ions is by collision with the cold atoms [6]. The ions are also subject to significant continuous rf heating as their numbers are large and $V_{IT} \gg V_{MOT}$, while the cooling is most effective only at the center of the MOT. Thus the ions are quite hot and therefore possess large velocities. Under such circumstances, even a glancing ion-atom collision, whether elastic or resonant charge exchange, will transfer sufficient energy to the laser-cooled atom such that it exceeds the capture velocity of the MOT and gets ejected from the MOT. This then forms the basis for equating the MOT loss rate with the ion-atom collision rate as described above.

Since the majority of ion-atom combination experiments use visible, laser-cooled ions, we provide a brief discussion of some major differences. In such cases, since ion temperatures are easily in the mK levels, the ion-atom cross sections are much higher. The ions then would be crystallized and have a velocity distribution that would be well characterized. How atoms interact with such crystallized ions and what part the

long-range order of the ions plays in the measurement of the binary ion-atom interaction needs to be carefully understood. A possible problem is that an ion-atom collision at these energies may not result in the ejection of the colliding atom from the MOT and would therefore require Eq. (18) to be written with a proportionality constant. However, if such an equation can be written, the cross section could be directly determined because of the well-characterized velocity and density distributions of the ions. Obviously, the rate equations constructed above would need to be modified according to the specifics of such a system.

VI. CONCLUSION

In a trapped ion and atom mixture, where the atoms are laser cooled and contained in a MOT and the ions are trapped within a Paul trap, we have developed a technique which measures the rate coefficient of ion-atom collisions, even when the ions are optically dark. The rate-equation formalism has been systematically developed for this experimental system. The experimental results are then fit to the rate-equation model, which is seen to provide a consistent and adequate description for the measurements. Analysis of the experimentally measured quantities with the model allows the determination of the trapped Rb^+ ions and Rb atom collision-rate coefficient. The value of the rate coefficient compares well with the theoretical estimation and validates the technique.

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