

Chapter 7

The gravitational lens system PKS 1830-211

7.1 Introduction

PKS 1830-211 is a compact flat-spectrum radio source and was identified as a candidate gravitational lens system by Rao & Subrahmanyan (1988) based on its peculiar morphology and spectral index. Its radio structure is made of two components, called the north-east (NE) and south-west (SW), separated by $\sim 1''$, each component consisting of a core-jet-knot configuration. Lower resolution observations (Jauncey, 1991) revealed an Einstein ring around the two components as well. Since this is the brightest known radio lens ($S_{1.4\text{GHz}}=10$ Jy), this is well studied in the cm and mm wavelengths. The components have comparable flux densities at cm-wavelengths and are highly variable at high frequencies ($\nu > 8$ GHz). Due to its proximity to the galactic centre ($l=12^\circ$, $b=-5.7^\circ$), the extinction in the optical as well as confusion in the infra-red makes the optical-IR identification of the background QSO and the lens galaxy difficult. Therefore the redshift of the background QSO was not known till the recent observations of Lidman (1999) when the NIR spectrum revealed the redshift to be 2.507 ± 0.002 . However, the redshift of the lens was known for quite some time, due to molecular absorption by the gas in the lens galaxy against the strong cm-mm continuum of the QSO. Wiklind & Combes (1996) discovered a host of molecular lines in absorption against the SW component, and derived the lens redshift to be 0.88582. Further observations (Wiklind & Combes, 1998) showed that there is much weaker molecular absorption against the NE component at the same redshift, but shifted in velocity by -147 km s^{-1} relative to the main absorption line against the SW image. HI and OH absorption were also detected by Chengalur, de Bruyn, & Narasimha (1999) using the WSRT. The strength of these absorption lines is maximum at the velocity of the molecular gas seen against the NE component and is weaker at the velocity corresponding to the SW component.

In addition to this lensing galaxy at $z=0.886$, Lovell et al. (1996) discovered a second lens system at $z=0.193$ through the detection of a feature at 1.2 GHz which they attributed to redshifted HI. Wiklind & Combes (1998) failed to detect any molecular line absorption at this redshift against either of the two components and derive an upper limit for the molecular gas mass of a few times $10^9 M_\odot$. Hence the reality of this system at $z=0.193$ has since remained in doubt. The modelling of the lens parameters of PKS 1830-211 is not unambiguous (see Nair, Narasimha, & Rao 1993 and Kochanek & Narayan 1992 for two different models) and hence

the possible existence of a second lens is important. Additionally, a measurement of the time delay in this lens has been obtained (time delay ~ 26 days; Lovell et al. 1998). Any attempt, therefore, at deriving the value of Hubble's constant from this lens system will depend on the detailed lens properties and geometry.

In order to confirm the existence of the second lens system at $z=0.193$, we searched for radio recombination line absorption at this redshift. This is a viable method of detecting this lens system since (1) An HI line has been seen and hence there exists substantial gas in this line of sight (2) the background source is strong and hence the geometry is favourable to detect stimulated recombination line emission due to the background continuum. Additionally, PKS 1830-211 is known to be highly variable at high frequencies and a single-epoch continuum spectrum would be useful for comparison. Hence we have also carried out observations to measure the single-epoch continuum flux density of PKS 1830-211 in the radio band.

7.2 Observations

We have carried out two different observations : (1) to measure the single-epoch continuum spectrum from 0.3–45 GHz using the VLA, (2) to search for RRLs at 1.4 GHz using the VLA. The observational procedure and the results for these observations are discussed below.

Table 7.1: Radio continuum spectrum of PKS 1830-211

Frequency (GHz)	Bandwidth (MHz)	Beam ("×")	rms (1σ) ^a (mJy)	Flux density (Jy)	Spectral ^c index
0.332	3.125	141×141	15	8.8±0.5	+0.09 ± 0.04
1.39	25	95×50	4	9.95±0.04	−0.14 ± 0.015
4.86	50	30×11	1	8.40±0.01	−0.03 ± 0.035
8.45	50	18×6	0.5	8.27±0.02	−0.43 ± 0.03
14.94	50	26×13	0.6	6.46±0.03	−0.68 ± 0.07
22.46	50	5×3	0.7	4.9±0.1	−1.1 ± 0.1
43.34 ^b	50	3×1.5	2	2.4±0.1	—

^aThe 1σ noise listed is for each IF.

^bThe 45 GHz observations were carried out using only 8 antennas.

^cSpectral index is calculated between that frequency and the one listed below.

7.2.1 Single-epoch radio continuum spectrum

The VLA was used in the D configuration on 27 October 1997 in the 4IF mode to measure the flux density of the source at all frequencies available at the VLA. The amplitudes were calibrated using 3C 286 at all frequencies with appropriate uv ranges. No phase calibrator was used. PKS 1830-211 is unresolved at all frequencies in this configuration and is bright even at the highest frequency. Hence the phases and amplitudes of the on-source visibilities were self-calibrated every 10 seconds, which was the integration time used. Final self-calibrated images were made for each IF at all the observed frequencies and flux densities were measured at each frequency. The 45 GHz observations were carried out without performing reference pointing. Further observational parameters and the measured flux densities are listed in Table 7.1. The 1σ rms in the single-IF images are also listed. The errors in the flux density measurements

are derived from the images; the true uncertainties are probably larger. The data are plotted in figure 7.1 with error-bars which are assumed to be the greater of the tabulated value and 1% of the continuum flux density.

Continuum spectrum of PKS 1830–211

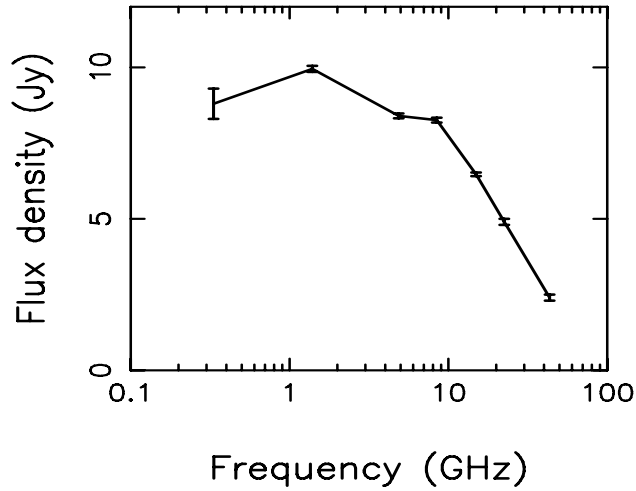


Figure 7.1: Single-epoch total flux density of PKS 1830-211 from 0.3 GHz to 45 GHz for the unresolved source. Error bars are the larger of the values listed in Table 7.1 and 1 % of the measured flux density. Observations were carried out on 27 October 1997.

7.2.2 Observations of H158 α and H136 α RRLs

The VLA was used in the A configuration at 1.4 GHz to search for RRLs in stimulated emission against the background radio continuum of the QSO. The correlator was used in the 4IF mode, yielding two different frequencies (IFs) of observation. One IF was tuned to detect the redshifted H158 α line from the $z=0.193$ system and the other was tuned to detect the redshifted H136 α from the $z=0.886$ system at the same time. These quantum numbers were chosen so that both the observing frequencies lie in the 1.4 GHz band of the VLA. Further observational parameters are listed in Table 7.2. A subsequent observation was made to detect the redshifted H158 α RRL, leading to a lower noise level for this line data. In neither case was a line detected. The upper limits to the line optical depth are listed in Table 7.2. Extremely sensitive upper limits could be derived since (1) the background source is strong ($S=10$ Jy) and (2) it is unresolved and hence accurate self-calibration could be performed.

Table 7.2: VLA search for RRLs at 1.4 GHz :
Observational Parameters

Parameter	$z=0.193$	$z=0.886$
Recombination line	H158 α	H136 α
Rest frequency (GHz)	1.652	2.586
Observed frequency (GHz)	1.384	1.371
Bandwidth (MHz)	1.56	3.125
Velocity resolution (km s $^{-1}$)	5.3	21
τ_{line} upper limit (5σ)	5×10^{-5}	5×10^{-4}

7.3 Modelling the ionized gas

We assume that the continuum spectrum of PKS 1830-211 at frequencies below 5 GHz is indeed flat (Rao & Subrahmanyan, 1988). Any observed deviation from flatness can then be attributed to free-free absorption by ionized gas in the intervening lens galaxies. From the measured flux densities at 1.4 GHz and 330 MHz, we derive an upper limit of 0.13 for the continuum optical depth at 330 MHz. We assume that the beam filling factor of any absorbing gas is unity. The (beam averaged) emission measure of the gas, for an electron temperature of 7500 K, is therefore $\leq 4 \times 10^4 \text{ pc cm}^{-6}$ if the gas is at $z=0.193$ and is $\leq 10^5 \text{ pc cm}^{-6}$ if the gas is at $z=0.886$.

The derived upper limits to the RRLs are modelled as follows. The ionized gas is assumed to have a uniform density n_e and a line-of-sight path length l pc. This gas is modelled as an uniform slab in front of the SW and NE components of the background QSO. Each component is assumed to have a 1.4 GHz flux density of 5 Jy and a spatial extent of $0.3''$. The values of n_e and l are varied between 10^{-2} – 10^6 cm^{-3} and 10^{-6} – 10^4 pc respectively. For each combination of n_e and l , the free-free continuum flux density and the RRL emission (arising due to spontaneous as well as stimulated emission) is calculated. Recently, Verheijen, Carilli, & Yun (2001) confirmed the existence of the HI line at $z=0.193$. Using the VLA along with the Pie-Town antenna, they were able to resolve the two components and detected the $z=0.193$ HI line against both the components with roughly equal optical depth, separated by $\sim 40 \text{ km s}^{-1}$. A line width of 80 km s^{-1} was assumed for calculating the expected line strengths. The model predictions were compared with the following two constraints: (1) the predicted line emission must be less than the observed upper limit and (2) the continuum emission from the ionized gas must be an insignificant fraction of the observed continuum at any frequency, i.e., the free-free opacity must not affect the continuum spectrum above 5 GHz significantly. For every assumed density, the maximum value of l and hence emission measure EM consistent with the above constraints was derived. This modelling scheme was carried out for both the redshift systems. Hence we obtain the maximum allowable (beam averaged) EM of the ionized gas as a function of the assumed density for both the $z=0.193$ and the $z=0.886$ lens galaxies and these are plotted in Figure 7.2.

This figure shows that our experiment is most sensitive to low density gas ($n_e = 1\text{--}10 \text{ cm}^{-3}$) and indicates an upper limit to the *beam averaged* emission measure of 100 pc cm^{-6} for this gas in the $z=0.193$ system. If the line emitting gas is predominantly of density $5\text{--}10 \text{ cm}^{-3}$, then the size of such a region located anywhere in the observed beam is constrained to be less than 200–300 pc in size, assuming a homogeneous gas distribution. On the other hand, if the gas is in compact structures with density $\sim 10^3 \text{ cm}^{-3}$, then its beam filling factor is constrained to be less than 6×10^{-4} . The limits to the $z=0.886$ system are about ten times higher and will not be discussed further.

7.4 Discussion

It is clear from the previous section that detecting classical HII regions from the lens system at $z=0.193$ might not be feasible with present day instruments. Due to the frequency dependence of the non-LTE optical depth, the experiment described above is most sensitive to low density gas. A similar observation at lower frequencies is desirable. RRL emission at 330 MHz, for

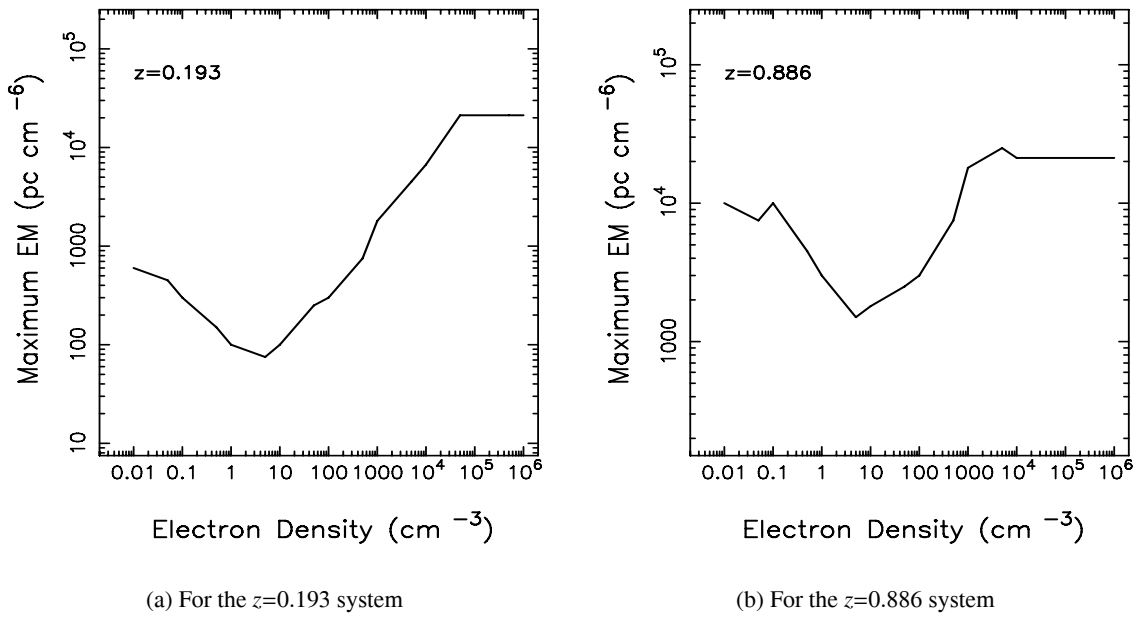


Figure 7.2: The maximum emission measure (EM) of the ionized gas in the **(a)**. $z=0.193$ system and the **(b)**. $z=0.886$ system which is consistent with the upper limits to the recombination line strengths, as a function of the assumed electron density. The kinks in the two curves are a result of numerical errors and are not significant.

example, will be more sensitive to the extended low density warm ionized medium (ELDWIM; Heiles, Reach, & Koo 1996; Roshni & Anantharamaiah 2001) in this galaxy. Though the primary motivation for the RRL observations was to confirm the existence of the $z=0.193$ system, it seems that a search for RRLs from the $z=0.886$ system might be more fruitful, since this system has a higher optical depth in various molecular lines. A 1.4 GHz RRL search at a sensitivity comparable to that achieved for the $z=0.193$ system as well as a search for lines at 330 MHz using the GMRT with a 1σ optical depth level of $\sim 10^{-5}$ should be attempted. The possibility of simultaneously observing multiple lines with the upgraded WSRT makes this optical depth limit attainable.

It is interesting that about two decades ago, it was expected that RRLs would be seen primarily from high redshift objects due to stimulated emission in the foreground ionized gas by the background radio bright nuclei (Shaver, 1978). In fact, study of RRLs at high redshifts were even thought to be useful for the determination of cosmological parameters (Sarazin & Wadiak, 1983) ! Hence, many searches had been conducted towards distant quasars (Churchwell & Shaver, 1979; Bell et al., 1984), but with no success. It was also generally believed that normal galaxies do not possess ionized gas of sufficient emission measure for RRLs produced purely by internal (or spontaneous) emission to be detectable. The observation described in this chapter seemed to offer the best possible chance to detect high redshift lines but were not fruitful. It is indeed ironic that all the extragalactic radio recombination line detections till date have been predominantly internal emission from nearby galaxies instead.

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