Singularities of the Mixed State Phase

In an interesting contribution to the physics of interference of particles with internal degrees of freedom, Sjöqvist et al. [1] have shown that if a beam of such particles, in a mixed state with a density matrix ρ_0 , is split in a Mach-Zhender interferometer and a unitary transformation U_i acts on the space of N internal states in one of the two paths, a phase difference given by $\arg Tr(U_i \rho_0)$ is introduced between the two beams. Here "phase difference" is defined as the shift in the maximum of the interference pattern. For pure states $|\psi_0\rangle$, this quantity reduces to the phase shift $\arg(\langle \psi_0 \mid U_i \mid \psi_0 \rangle)$ defined by Pancharatnam [2]. The purpose of this Comment is to point out that the definition of the mixed state phase as in Ref. [1] breaks down at points in the parameter space where $|\operatorname{Tr}(U_i\rho_0)| = 0$. Equation (8) in the paper shows that at such points the interference pattern has zero contrast. Such points constitute singularities of phase involving discontinuous phase jumps. In the pure state case, such singularities have been demonstrated theoretically and experimentally for N = 2 using the two polarization states of light [3-6]. The mixed state phase involves singularities in new kinds of parameter spaces which include variables representing decoherence of quantum states. We show that the singularities make interpretation of even simple interference experiments like the ones analyzed in Ref. [1], in terms of the proposed phase, nontrivial.

Consider a beam of spin-1/2 particles in a mixed state with a density matrix $\rho_0 = \text{diag}[(1 + r)/2, (1 - r)/2]$ in the basis of $|z\pm\rangle$; $|z\pm\rangle$ being the eigenstates of z component of the spin and r the degree of polarization. The beam is split symmetrically in an interferometer and, to keep things simple, a unitary transformation U_i , diagonal in the $|z\pm\rangle$ basis, is applied in one of the two paths by means of a variable magnetic field along \hat{z} . This results in the $|z\pm\rangle$ states acquiring a phase factor $e^{\pm i\delta}$. The mixed state phase ϕ is given by $Tr(U_i \rho_0) = [\cos \delta + ir \sin \delta] =$ $ce^{i\phi}$. In the plane of the parameters (r, δ) , phase singularities occur at the points $(0, (i + 1/2)\pi)$, where c vanishes; *j* being any integer. The interference pattern for the unpolarized case (r = 0) can be looked upon as a superposition of two cosine intensity patterns with equal amplitude for the two spin eigenstates, moving in opposite directions as δ is varied [5]. When $\delta = (i + 1/2)\pi$, the two patterns have a relative phase shift $(2i + 1)\pi$, yielding uniform illumination, making the phase indeterminate. A closed circuit in the (r, δ) plane around any one of these singularities yields a value $\pm 2\pi$ for the total phase change $\int d\phi$.

Figure 1 shows a set of typical phase shifts acquired by a mixed state as a function of δ . In each case the phase shift equals π in magnitude for a variation of δ through π , is linear for pure states, i.e., for $r = \pm 1$ [curves 1(*A*) and 1(*B*)], is highly nonlinear near the singularity at



FIG. 1. Computed phase shift $\int d\phi$ as a function of applied magnetic field for (A) r = 1, (B) r = -1, (C) r = 0.001, and (D) r = -0.001.

 $(r = 0, \delta = \pi/2)$ [curves 1(C) and 1(D)], and changes sign with the sign of $r (d\phi/d\delta)$ has the same sign as r). For r = 0, the phase shift is indeterminate and not π as claimed in [1]. The unpolarized (r = 0) neutron interferometer experiments cited therein do not measure the phase shift curves as shown in Fig. 1. These can, however, be measured in an interference experiment with neutrons or with polarized light if r can be varied and phase shifts measured in the experiment. The basic topological effect, namely, a $2n\pi$ phase change (*n* being integer), resulting from a circuit around one or more singularities can in fact be measured without a close approach to the singularities, making experiments easier. In an experiment with spin > 1/2, even a circuit around a single singularity could result in a $2n\pi$ phase shift with n > 1, yielding information about the spin quantum number. A nonmodular description of the phase, as in Refs. [3,4,5] thus addresses an essential aspect of the topological phase, missed in a modulo 2π description [1].

Rajendra Bhandari^{*} Raman Research Institute, Bangalore 560080, India

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*Electronic address: bhandari@rri.res.in

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